Planar nanocontacts with atomically controlled separation

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We have developed a technology to reproducibly make gaps with distance control on the single-atom scale. The gold contacts are flat on the nanometer scale and are fabricated on an oxidized aluminum film that serves as a gate. We show that these contacts are clean and can be stabilized via chemical functionalization. Deposition of conjugated molecules leads to an increase in the gap conductance of several orders of magnitude. Stable current–voltage characteristics at room temperature are slightly nonlinear. At low temperature, they are highly nonlinear and show a clear gate effect. © 2003 American Institute of Physics. [DOI: 10.1063/1.1623317]

An important problem in molecular electronics consists of fabricating three-terminal devices with nano-objects. Several methods have been developed.1 In the mechanical approach (scanning tunneling microscopy and break junctions2,3) piezoelectric control of the gap distance offers great precision well below the atomic scale. Transport through small molecules have been measured,2–6 but implementation of a gate is difficult. Electromigration7 of gold wires on top of aluminum or silicon gate is in that respect more promising and has shown impressive results.8,9 Nevertheless, the electromigration method only provides statistical etching of very thin gold films (separation ranging from one atom up to a few nanometers). Is then a problem if the smallest distance is far from the substrate.

In this letter, we show that flat, stable contacts with a separation ranging from one atom up to a few nanometers can be made on aluminum gates. We use electrochemical etching of very thin gold films (~15 nm) while monitoring the conductance of the gap. Subatomic resolution is obtained with a high yield. After rinsing in demineralized water, the contacts are clean enough to be directly functionalized with molecules containing thiol or sulfide groups. Measurements of $I–V$ characteristics show that they then form stable contacts for nano-objects.

Fabrication consists of two major steps: evaporation of a thin gold link on top of an oxidized aluminum film and, secondly, electrochemical etching of the gold to form a nanogap. In the first major step, we start with evaporating a 20-nm-thick, 600-nm-wide aluminum wire on an oxidized undoped silicon substrate (see Fig. 1). On top, we define two lines (the electrodes) separated by about 100 nm using e-beam lithography in a two-layer resist. After resist development, the evaporation sequence is as follows: first a 7 nm sticking layer of titanium is evaporated perpendicularly. The sample is then tilted by 15° and the 15 nm thin gold link is formed. Finally, the thick gold pads are deposited at 90°.

In the second major step, we slowly reduce the weak link thickness using electrochemical etching in a bath of KAuCN$_2$ (0.1 mM of KAuCN$_2$, 0.2 M of KOH, and 1 M of KHCO$_3$). As counterelectrode, we use a small piece of silicon (~20 mm$^2$) covered with 100 nm evaporated gold. The voltage on this counterelectrode ($V_{CE}$) sets the etching rate. The electrodes are contacted using spring contacts, and only a small part of the gold electrodes (~0.05 mm$^2$) is dipped in the etching bath. During the etching process, the interelectrode conductance measured using a $V–I$ converter and a lock-in amplifier to detect the small ac responses. Gap resistance can be measured up to 1 GΩ. Note that from comparing the gap resistance in solution and after drying, we conclude that conduction enhancement by the solution is not significant.

The electrochemical etching proceeds in two parts. We start etching the gold at a high current density with $V_{CE} = -2400$ mV. The excitation frequency of the lock-in is set to 41 Hz13 and the measured signal is fed through a high pass filter to remove the large dc etching current. After a few minutes, the interelectrode conductance starts to drop significantly. When the resistance reaches a value of about 3 kΩ, $V_{CE}$ is changed to $-1560$ mV. The etching rate is now so slow that removal of a single atom from the gap region can be observed. Figure 2 shows one out of the 100 very similar

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traces of the current through the gap versus etching time. We find that when the conductance is around $G_0 = 2e^2/h$, the current trace becomes more noisy for a time span of about $t_0$. A sharp decrease in the current follows and the gap opening starts to show large fluctuations. These observations suggest that $t_0$ is the typical time scale to remove the last gold atom from the contact. The noise that occurs after opening of the gap remains after washing and drying. We attribute it to gold-atom motion on the surface. More details of the fabrication process and the stability will be published elsewhere.\(^{14}\)

An interesting feature of our data is that the etching rate is reproducible even at the atomic scale. As shown in Fig. 2, the gap resistance increases close to exponential with etching time after opening of the gap. Since tunneling is exponential in distance, this indicates an approximate constant etching rate on the atomic scale. Moreover for all our samples, we find that the time span to open the gap to $1 \, \text{G}\Omega$ equals $\sim 3 \, t_0$ when starting from the sudden decrease of the current. A gap with a resistance $1 \, \text{G}\Omega$ is estimated\(^{15,16}\) to be 3 atoms wide ($\sim 1 \, \text{nm}$), in agreement with our previous estimate that $t_0$ is the time scale to remove one gold atom. We also find that one time $t_0$ after the last atom removal, the resistance is always close to $1 \, \text{M}\Omega$. This value then corresponds to a gap of one atom wide. Furthermore, as illustrated in Fig. 2, the noise has an amplitude that corresponds to the motion of gold atoms over half an atomic distance ($\sim 1.5 \, \text{Å}$), indicating that atomic resolution is achieved.

After rinsing in water and drying the gaps, we have checked the gate at room temperature. We generally find very small leakage currents (smaller than $10 \, \text{pA}$ at $1 \, \text{V}$ for more than $50\%$ of the samples). The typical gate voltage beyond which current starts to grow significantly is $\pm 2 \, \text{V}$. Importantly, scanning electron microscope (SEM) imaging also indicates that the aluminum gates are not affected by the etching process.

We have studied the morphology of the thin etched electrode with a SEM. Near the gap, we find grains that vary in size between 15 and 8 nm. Smaller grains are not observed even at high resolution. The gap region is generally barely visible, indicating an extreme shallowness compared to other remaining parts of the thin gold link (with atomic force microscopy we were also unable to resolve the height of the end parts of the electrodes). This clearly indicates that the geometry is planar and that the point of smallest gap distance is directly on top of the aluminum gate. The use of aluminum is crucial: thin gold films on silicon oxide break during etching before atomic resolution is obtained. Experimentally and theoretically it has been shown that in ultrathin films on aluminum oxide, gold atoms strongly bind to the negatively charged oxygen sites, which may account for the greater stability.\(^{15,16}\)

By measuring the gap resistance, we find that bare gaps are stable on the scale of about one hour in ambient conditions. Stability is drastically improved by functionalization of the surface by molecules with thiol or sulfide ends. As a test that the gaps can be used for transport measurements of single or a small number of molecules, we have dipped them in millimolar ethanolic solutions of 1-3 disulfobenzene. The increase in current is of more than two orders of magnitude.\(^{17,18}\) Inset: $I-V$ curves at $4.2 \, \text{K}$ for two values of the gate voltage. As can clearly be seen, the gap of $\sim 300 \, \text{mV}$ can completely be suppressed. For all samples the deposition time was $24 \, \text{h}$.

![FIG. 2. Gap current (upper graph) and resistance (lower graph) versus etching time. Only the slow etching part, in which the gap is formed, is shown. $V_{CE}$ is kept constant at $-1470 \, \text{mV}$. The excitation bias was $10 \, \text{mV}$ with a frequency of $5 \, \text{Hz}$. For $t > 2000 \, \text{s}$, the gap resistance shows an approximately exponential time dependence. The dashed lines correspond to a resistance deviation equivalent to about 0.5 atomic length. Time indication \(t_0\) has been added to characterize the last plateau length and the gap opening. The current value $2e^2/h \, \text{V}$ indicates the conductance quantum.](image1)

![FIG. 3. (a) Room-temperature $I-V$ characteristic for a four-atom gap covered with 1-3 disulfobenzene. The increase in current is of more than two orders of magnitude. (b) Characteristics for a gap covered with oligo(cyclohexylidenes). Inset: $I-V$ curves at $4.2 \, \text{K}$ for two values of the gate voltage. As can clearly be seen, the gap of $\sim 300 \, \text{mV}$ can completely be suppressed. For all samples the deposition time was 24 h.](image2)
atoms wide (to be more precise corresponding to 4–5 $t_0$) show a resistance of $\sim 1 \text{ M} \Omega$. For slightly larger gaps corresponding to 6–8 $t_0$, resistances of the order of 10 M$\Omega$ are reproducibly found.

Several experiments have been performed to show that transport through the molecules is responsible for the strongly enhanced current. First, we have dipped over 10 gaps in ethanol for more than 24 h and did not find any decrease of the resistance. In addition, for gaps that are larger than 2.5 molecular length (so that two molecules facing each other but attached to different electrodes do not touch), the resistance after dipping in the oligo(cyclo hexylidenes) solution is orders of magnitude larger (i.e., $> 1 \text{ G}\Omega$).

$I–V$ traces have been measured over period of ten hours at room temperature, and we find that the stability is much better than bare gaps. Occasionally, sudden switchings occur, indicating reconfigurations of the gold atoms contacting the molecules. Furthermore, when kept in vacuum the resistance of the molecular nanodevice remains approximately the same (over the time span of weeks).

Preliminary measurements have been performed at low temperatures. $I–V$ characteristics show gap-like curves, as illustrated in the inset of Fig. 3. For the four samples studied, the same gap of $\sim 300$ mV was found; for the three samples with the gate connected, a clear gate effect has been observed (see inset of Fig. 3). Interestingly, the gate effect presents unconventional features (no saturation, nor periodicity as observed in Coulomb blockade). A systematic study of the gate effect and the low temperature behavior is beyond the scope of this letter and will be presented elsewhere.

In conclusion, we have made flat gold nanogaps with atomic precision. The fabrication process is compatible with aluminum gating technology and yields gaps that can readily be functionalized with standard dithiol linkers. Once bound chemically to rigid nano-objects, these gaps are stable for weeks. Our method is compatible with current on-chip technology and provides a simple, fast, and reproducible technique for the fabrication of molecular and few atoms based nanodevices.

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12. To perform the maximum resolution ($\sim 1 \text{ pA}$), the frequency has to be reduced to a few hertz at low etching current to minimize the influence of parasitic capacitances.